

Topological nature of the Fu-Kane-Mele invariants

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1. *Time reversal symmetries and “Quaternionic” structures*
2. *The role of the (involutive) base space*
3. *In the search of a classifying object*
4. *FKMM vs. Fu-Kane-Mele*

Topological Quantum Systems with odd TRS's

Let \mathbb{B} a topological space, (“**Brillouin zone**”). Assume that:

- \mathbb{B} is a CW-complex (compact, Hausdorff and path-connected);

DEFINITION (Topological Quantum System (TQS))

Let \mathcal{H} be a separable Hilbert space and $\mathcal{K}(\mathcal{H})$ the algebra of compact operators. A **TQS** is a self-adjoint map

$$\mathbb{B} \ni k \mapsto H(k) = H(k)^* \in \mathcal{K}(\mathcal{H})$$

continuous with respect to the norm-topology.

👉 The **spectrum** $\sigma(H(k)) = \{E_j(k) \mid j \in \mathcal{I} \subseteq \mathbb{Z}\} \subset \mathbb{R}$, is a sequence of eigenvalues ordered according to

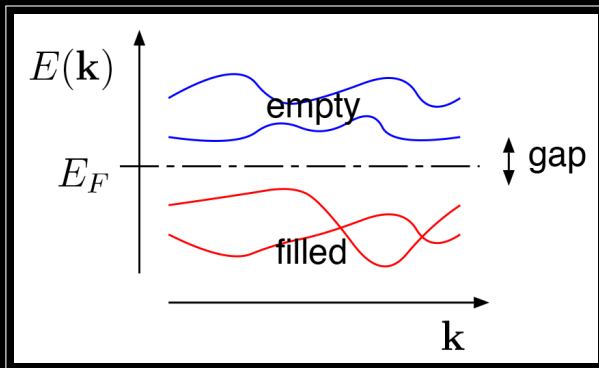
$$\dots E_{-2}(k) \leq E_{-1}(k) < 0 \leq E_1(k) \leq E_2(k) \leq \dots$$

👉 The maps $k \mapsto E_j(k)$ are **continuous** (energy bands) ...

Topological Quantum Systems with odd TRS's

... namely a band spectrum

$$H(\mathbf{k}) \psi_j(\mathbf{k}) = E_j(\mathbf{k}) \psi_j(\mathbf{k}), \quad \mathbf{k} \in \mathbb{B}$$



Usually an energy **gap** separates the filled **valence** bands from the empty **conduction** bands. The **Fermi level** E_F characterizes the gap.

Topological Quantum Systems with odd TRS's

A homeomorphism $\tau : \mathbb{B} \rightarrow \mathbb{B}$ is called **involution** if $\tau^2 = \text{Id}_{\mathbb{B}}$. The pair (\mathbb{B}, τ) is called an **involutive space** and $\mathbb{B}^{\tau} \subset \mathbb{B}$ is the subsetset of **invariant points**. Each space \mathbb{B} admits (at least) the **trivial involution** $\tau_{\text{triv}} := \text{Id}_{\mathbb{B}}$.

DEFINITION (TQS with time-reversal symmetry)

Let (\mathbb{B}, τ) be an involutive space, \mathcal{H} a separable Hilbert space endowed with a **complex conjugation** \mathbf{C} . A TQS $\mathbb{B} \ni \mathbf{k} \mapsto H(\mathbf{k})$ has a **time-reversal symmetry** (TRS) of parity $\eta \in \{\pm 1\}$ if there is a continuous unitary-valued map $\mathbf{k} \mapsto U(\mathbf{k})$ such that

$$U(\mathbf{k}) H(\mathbf{k}) U(\mathbf{k})^* = \mathbf{C} H(\tau(\mathbf{k})) \mathbf{C}, \quad \mathbf{C} U(\tau(\mathbf{k})) \mathbf{C} = \eta U(\mathbf{k})^* .$$

A TQS with an **odd** TRS (i.e. $\eta = -1$) is called of class **AII**.

The Serre-Swan construction

- An **isolated family** of energy bands is any (finite) collection $\{E_{j_1}(\cdot), \dots, E_{j_m}(\cdot)\}$ of energy bands such that

$$\min_{k \in \mathbb{B}} \text{dist} \left(\bigcup_{s=1}^m \{E_{j_s}(k)\}, \bigcup_{j \in \mathcal{I} \setminus \{j_1, \dots, j_m\}} \{E_j(k)\} \right) = C_g > 0.$$

This is usually called **gap condition**.

- An isolated family is described by the **Fermi projection**

$$P_F(k) := \sum_{s=1}^m |\psi_{j_s}(k)\rangle \langle \psi_{j_s}(k)|.$$

This is a **continuous** projection-valued map

$$\mathbb{B} \ni k \mapsto P_F(k) \in \mathcal{K}(\mathcal{H}).$$

The Serre-Swan construction

☞ For each $k \in \mathbb{B}$

$$\mathcal{H}_k := \text{Ran } P_F(k) \subset \mathcal{H}$$

is a subspace of \mathcal{H} of **fixed** dimension m .

☞ The collection

$$\mathcal{E}_F := \bigsqcup_{k \in \mathbb{B}} \mathcal{H}_k$$

is a topological space (said **total space**) and the **map**

$$\pi : \mathcal{E}_F \longrightarrow \mathbb{B}$$

defined by $\pi(k, v) = k$ is continuous (and open).

This is a (rank- m) **complex** vector bundle called **Bloch-bundle**.

The Serre-Swan construction

👉 An **odd** TRS induces a **“Quaternionic”** structure on the Bloch-bundle.

DEFINITION (Atiyah, 1966 - Dupont, 1969)

Let (\mathbb{B}, τ) be an involutive space and $\mathcal{E} \rightarrow \mathbb{B}$ a **complex** vector bundle. Let $\Theta : \mathcal{E} \rightarrow \mathcal{E}$ an **homeomorphism** such that

$$\Theta : \mathcal{E}|_k \longrightarrow \mathcal{E}|_{\tau(k)} \quad \text{is } \mathbf{anti}\text{-linear} .$$

[\mathcal{R}] - The pair (\mathcal{E}, Θ) is a **“Real”**-bundle over (\mathbb{B}, τ) if

$$\Theta^2 : \mathcal{E}|_k \xrightarrow{+1} \mathcal{E}|_k \quad \forall k \in \mathbb{B} ;$$

[\mathcal{Q}] - The pair (\mathcal{E}, Θ) is a **“Quaternionic”**-bundle over (\mathbb{B}, τ) if

$$\Theta^2 : \mathcal{E}|_k \xrightarrow{-1} \mathcal{E}|_k \quad \forall k \in \mathbb{B} .$$

The classification problem

DEFINITION (Topological phases)

Let $\mathbb{B} \ni k \mapsto H(k)$ be an **odd TR-symmetric** TQS with an isolated family of m energy bands and associated “**Quaternionic**” Bloch bundle $\mathcal{E}_F \rightarrow \mathbb{B}$. The **topological phase** of the system is specified by

$$[(\mathcal{E}_F, \Theta)] \in \text{Vec}_Q^m(\mathbb{B}, \tau).$$



Main Question:

How to classify $\text{Vec}_Q^m(\mathbb{B}, \tau)$ at least for **low-dimensional** \mathbb{B} ?

The classification problem

Known results for $\dim(\mathbb{B}) \leq 3$

- $\text{Vec}_{\mathbb{C}}^m(\mathbb{B}) \stackrel{c_1}{\simeq} H^2(\mathbb{B}, \mathbb{Z})$ (Peterson, 1959)
- $\text{Vec}_{\mathbb{R}}^m(\mathbb{B}, \tau) \stackrel{c_1^{\mathbb{R}}}{\simeq} H_{\mathbb{Z}_2}^2(\mathbb{B}, \mathbb{Z}(1))$ (Kahn, 1987 - D. & Gomi, 2014)

<i>CAZ</i>	<i>TRS</i>	<i>Category</i>	<i>VB</i>
<i>A</i>	0	<i>complex</i>	$\text{Vec}_{\mathbb{C}}^m(\mathbb{B})$
<i>AI</i>	+	<i>"Real"</i>	$\text{Vec}_{\mathbb{R}}^m(\mathbb{B}, \tau)$
<i>AII</i>	-	<i>"Quaternionic"</i>	$\text{Vec}_{\mathbb{Q}}^m(\mathbb{B}, \tau)$

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Electrons in a periodic environment

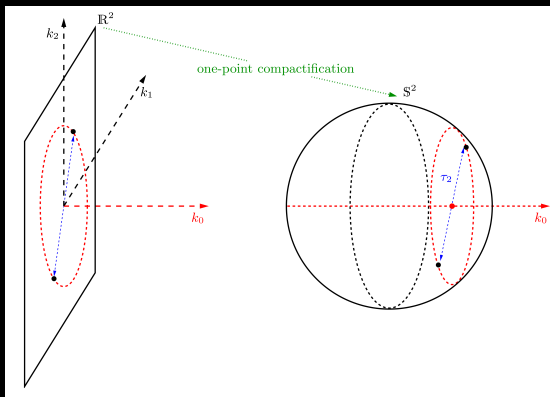
- **Periodic** quantum systems (e.g. absence of **disorder**):
 - \mathbb{R}^d -translations \Rightarrow **free** (Dirac) fermions;
 - \mathbb{Z}^d -translations \Rightarrow **crystal** (Bloch) fermions.
- The Bloch-Floquet (or Fourier) theory exploits the invariance under translations of a periodic structure to describe the state of the system in terms of the **quasi-momentum** \mathbf{k} on the **Brillouin zone** \mathbb{B} .
- Complex conjugation (TRS) endows \mathbb{B} with an involution τ .
- Examples are:
 - Gapped electronic systems,
 - BdG superconductors,
 - **Photonic crystals** (M. Lein talk).

Continuous case $\mathbb{B} \equiv \mathbb{S}^{1,d}$

$$\mathbb{S}^d \xrightarrow{\theta_{1,d}} \mathbb{S}^d$$

$$\mathbb{S}^{1,d} := (\mathbb{S}^d, \theta_{1,d})$$

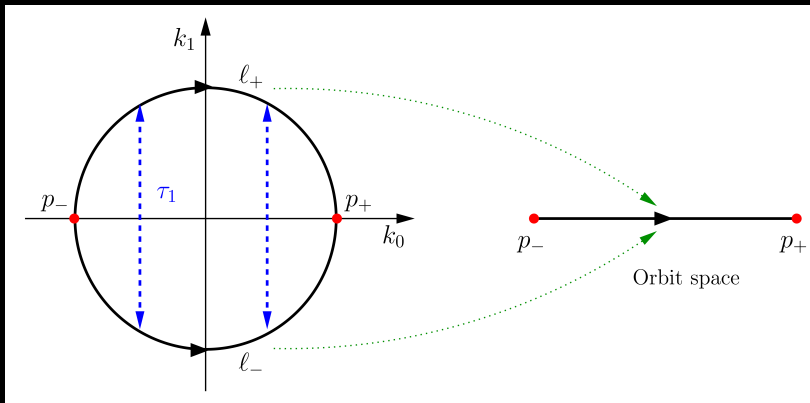
$$(+k_0, +k_1, \dots, +k_d) \xrightarrow{\theta_{1,d}} (+k_0, -k_1, \dots, -k_d)$$



Periodic case $\mathbb{B} \equiv \mathbb{T}^{0,d,0}$


$$\mathbb{S}^{1,1} \times \dots \times \mathbb{S}^{1,1} \xrightarrow{\tau_d := \theta_{1,1} \times \dots \times \theta_{1,1}} \mathbb{S}^{1,1} \times \dots \times \mathbb{S}^{1,1}$$

$$\mathbb{T}^{0,d,0} := \underbrace{\mathbb{S}^{1,1} \times \dots \times \mathbb{S}^{1,1}}_{d \text{ - times}} = (\mathbb{T}^d, \tau_d)$$



Topological states for Bloch electrons

	$d = 1$	$d = 2$	$d = 3$	$d = 4$	
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{S}^{1,d})$	0	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}	Free
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{T}^{0,d,0})$	0	\mathbb{Z}_2	\mathbb{Z}_2^4	$\mathbb{Z}_2^{10} \oplus \mathbb{Z}$	Periodic

 First proof for $d = 1, 2$ due to **Fu, Kane and Mele (2005 - 2007)** based on the

$$\text{Fu-Kane-Mele index} := \prod_{k_i \in \mathbb{B}^T} \frac{\sqrt{\det[\mathbf{W}(k_i)]}}{\text{Pf}[\mathbf{W}(k_i)]}.$$

Here $\mathbb{B}^T \ni \mathbf{k} \mapsto \mathbf{W}(\mathbf{k})$ is an antisymmetric matrix built from the Bloch functions.

!! It makes sense **only** when \mathbb{B}^T is **finite** !!

Topological states for Bloch electrons

An afterwards ...

- ☞ Computed by **Kitaev (2009)** for all d by **K-theory** (stable range).
- ☞ “Handmade” **frame construction** for the case $\mathbb{T}^{0,2,0}$ by **Graf and Porta (2013)** and for the case $\mathbb{T}^{0,3,0}$ by **Fiorenza, Monaco and Panati (2016)** and **Cornean, Monaco and Teufel (2016)**.
- ☞ **D. and Schulz-Baldes (2015)** with **spectral flux** (disorder).
- ☞ **Kennedy and Zirnbauer (2015)** by the calculation of the **equivariant homotopy** (very general but **hard** to compute).
- ☞ **D. and Gomi (2015)** by the introduction of the **FKMM-invariant** (a characteristic class) and the computation of the **equivariant cohomology** (very general and **not so hard** to compute).

Why more general involutive spaces?

☞ *The external triggering:*

\mathbb{B} can be interpreted as the space of **control parameters** for a quantum system **adiabatically perturbed**.

☞ *The Born-Oppenheimer approximation:*

Many systems depend by **slow** and **fast** degrees of freedom (e.g. the Molecular Dynamics). Under certain conditions the slow and fast variables **decouple adiabatically** (i.e. the fast variables **adjust instantly** to changes of the slow variables). As a consequence, the fast dynamics is described by an **effective Hamiltonian** which depends by the slow (**classical**) degrees of freedom. “De facto” one is in a situation described by a **TQS**

$$X \ni (q, p) \longmapsto H_{\text{fast}}(q, p)$$

with X the **classical phase space**. The **TR symmetry** acts on the classical variables and induces an involution on the space X .

Why more general involutive spaces?

Therefore (\mathbb{B}, τ) can be very general. In particular the **fixed-point set** \mathbb{B}^τ could be empty (**free** action) or a sub-manifold of whatever co-dimension (and not necessary a discrete set of points).

For instance there are family of involutive spheres $\mathbb{S}^{p,q} := (\mathbb{S}^{p+q-1}, \theta_{p,q})$ with $\theta_{p,q}$ defined by

$$(k_0, k_1, \dots, k_{p-1}, k_p, \dots, k_{p+q-1}) \xrightarrow{\theta_{p,q}} (k_0, k_1, \dots, k_{p-1}, -k_p, \dots, -k_{p+q-1})$$

and of involutive tori

$$\mathbb{T}^{a,b,c} := \underbrace{\mathbb{S}^{2,0} \times \dots \times \mathbb{S}^{2,0}}_{a\text{-times}} \times \underbrace{\mathbb{S}^{1,1} \times \dots \times \mathbb{S}^{1,1}}_{b\text{-times}} \times \underbrace{\mathbb{S}^{0,2} \times \dots \times \mathbb{S}^{0,2}}_{c\text{-times}}$$

Recently **Gat** and **Robbins** ([arXiv:1511.08994](https://arxiv.org/abs/1511.08994)) considered the cases $\mathbb{B} = \mathbb{S}^{0,3}$ (**rigid rotor**) and $\mathbb{B} = \mathbb{T}^{1,1,0}$ (**phase space of slow dynamic of a 1D periodic particle**). In the first case $\mathbb{B}^\tau = \emptyset$ and in the second $\mathbb{B}^\tau = \mathbb{S}^1 \sqcup \mathbb{S}^1$.

Why more general involutive spaces?

!! Many of the previous approaches just fail when \mathbb{B}^T is **not** a finite set !!



Which object replaces the **Fu-Kane-Mele index** when \mathbb{B}^T is not a finite set ?



A characteristic (cohomological) class called **FKMM-invariant**.

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Relative equivariant cohomology

The Borel's construction

- (X, τ) any involutive space and $(\mathbb{S}^\infty, \theta)$ the **infinite** sphere (contractible space) with the **antipodal** (free) involution:

$$X_{\sim\tau} := \frac{\mathbb{S}^\infty \times X}{\theta \times \tau} \quad (\text{homotopy quotient}).$$

- \mathcal{Z} any abelian ring (module, system of coefficients, ...)

$$H_{\mathbb{Z}_2}^j(X, \mathcal{Z}) := H^j(X_{\sim\tau}, \mathcal{Z}) \quad (\text{eq. cohomology groups}).$$

- $\mathbb{Z}(m)$ the **\mathbb{Z}_2 -local system** on X based on the module \mathbb{Z}

$$\mathbb{Z}(m) \simeq X \times \mathbb{Z} \quad \text{endowed with} \quad (x, \ell) \mapsto (\tau(x), (-1)^m \ell).$$


Relative equivariant cohomology

- $H_{\mathbb{Z}_2}^\bullet$ is a (generalized) **cohomology theory** which can be extended to pairs of spaces $Y \subseteq X$ in order to define **relative cohomology groups** $H_{\mathbb{Z}_2}^\bullet(X|Y, \mathbb{Z})$.
- In [D. - Gomi, 2015] we showed that

$$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{T}^{0,d,0}) \quad \text{and} \quad \text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{S}^{1,d-1}), \quad d \leq 4$$

can be classified by a **characteristic class** with values in $H_{\mathbb{Z}_2}^2(\mathbb{B}|\mathbb{B}^\tau, \mathbb{Z}(1))$: the **FKMM-invariant**.

$$\begin{array}{ccccc} H_{\mathbb{Z}_2}^1(\mathbb{B}^\tau, \mathbb{Z}(1)) & \xrightarrow{\delta_1} & H_{\mathbb{Z}_2}^2(\mathbb{B}|\mathbb{B}^\tau, \mathbb{Z}(1)) & \xrightarrow{\delta_2} & H_{\mathbb{Z}_2}^2(\mathbb{B}, \mathbb{Z}(1)) & \xrightarrow{r} & H_{\mathbb{Z}_2}^2(\mathbb{B}^\tau, \mathbb{Z}(1)) \\ \text{R} & & & & \text{R} & & \text{R} \\ [\mathbb{B}^\tau, \mathbb{S}^{1,1}]_{\mathbb{Z}_2} & & \text{Pic}_{\mathcal{R}}(\mathbb{B}, \tau) & & \text{Pic}_{\mathbb{R}}(\mathbb{B}^\tau) & & \end{array}$$

 The results in [D. - Gomi, 2015] only apply to the case of a \mathbb{B}^τ finite. To consider more general involutive spaces we need more generality!

The (generalized) FKMM-invariant

THEOREM (D. - Gomi, 2016 | Part I)

Given (\mathbb{B}, τ) let

$$\mathrm{Pic}_{\mathcal{R}}(\mathbb{B}|\mathbb{B}^{\tau}, \tau) := \{[(\mathcal{L}, \mathbf{s})] \mid \mathcal{L} \in \mathrm{Pic}_{\mathcal{R}}(\mathbb{B}, \tau), \mathbf{s} : \mathcal{L}|_{\mathbb{B}^{\tau}} \rightarrow \mathbf{U}(1)\}$$

with group structure given by the **tensor product**. Then

$$\mathrm{Pic}_{\mathcal{R}}(\mathbb{B}|\mathbb{B}^{\tau}, \tau) \stackrel{\tilde{\kappa}}{\simeq} H_{\mathbb{Z}_2}^2(\mathbb{B}|\mathbb{B}^{\tau}, \mathbb{Z}(1)) .$$

This result extends the **Kahn's isomorphism**

$$\mathrm{Vec}_{\mathcal{R}}^m(\mathbb{B}, \tau) \simeq H_{\mathbb{Z}_2}^2(\mathbb{B}, \mathbb{Z}(1))$$

and indeed can be proved in a similar way.

The (generalized) FKMM-invariant

THEOREM (D. - Gomi, 2016 | Part II)

There is a group homomorphism

$$\kappa : \text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{B}, \tau) \longrightarrow H_{\mathbb{Z}_2}^2(\mathbb{B}|\mathbb{B}^\tau, \mathbb{Z}(1))$$

called the **FKMM-invariant**.

(1) Determinant functor:

If $(\mathcal{E}, \Theta) \in \text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{B}, \tau)$ then $(\det \mathcal{E}, \det \Theta) \in \text{Pic}_{\mathcal{R}}(\mathbb{B}, \tau)$;

(2) Canonical section:

It exists a **unique** (canonical) trivialization

$$h_{\text{can}} : \det \mathcal{E}|_{\mathbb{B}^\tau} \rightarrow \mathbb{B}^\tau \times \mathbb{C}$$

which define $\mathbf{s}_{\text{can}}(k) := h_{\text{can}}^{-1}(k, 1)$ for all $k \in \mathbb{B}^\tau$.

(3) The mapping $\mathcal{E} \mapsto (\det \mathcal{E}, \mathbf{s}_{\text{can}}) \in \text{Pic}_{\mathcal{R}}(\mathbb{B}|\mathbb{B}^\tau, \tau)$:

$$\kappa(\mathcal{E}, \Theta) := \tilde{\kappa}(\det \mathcal{E}, \mathbf{s}_{\mathcal{E}}).$$

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Properties of $\kappa : \text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{B}, \tau) \longrightarrow H_{\mathbb{Z}_2}^2(\mathbb{B}|\mathbb{B}^\tau, \mathbb{Z}(1))$

- (i) **Isomorphic** \mathbb{Q} -bundles have the same FKMM-invariant;
- (ii) If (\mathcal{E}, Θ) is \mathbb{Q} -trivial then $\kappa(\mathcal{E}, \Theta) = \mathbf{0}$;
- (iii) κ is **natural** under the pullback induced by equivariant maps;
- (iv) $\kappa(\mathcal{E}_1 \oplus \mathcal{E}_2, \Theta_1 \oplus \Theta_2) = \kappa(\mathcal{E}_1, \Theta_1) + \kappa(\mathcal{E}_2, \Theta_2)$;
- (v) κ is the image of a **universal class** $\mathfrak{h}_{\text{univ}}$;
- (vi) When $\mathbb{B}^\tau = \{\text{finite collection of points}\}$ then
$$\kappa(\mathcal{E}, \Theta) \simeq \text{Fu-Kane-Mele invariants};$$
- (vii) When $\mathbb{B}^\tau = \emptyset$
$$\kappa(\mathcal{E}, \Theta) \simeq c_1^{\mathcal{R}}(\det \mathcal{E}, \det \Theta);$$
- (viii) When $\mathbb{B}^\tau = \emptyset$ and $\text{Pic}_{\mathbb{Q}}(\mathbb{B}, \tau) \neq \emptyset$ then $\text{Pic}_{\mathbb{Q}}(\mathbb{B}, \tau)$ is a **torsor** over $\text{Pic}_{\mathcal{R}}(\mathbb{B}, \tau)$. Hence

$$\text{Pic}_{\mathbb{Q}}(\mathbb{B}, \tau) \simeq \text{Pic}_{\mathcal{R}}(\mathbb{B}, \tau) \simeq H_{\mathbb{Z}_2}^2(\mathbb{B}, \mathbb{Z}(1)).$$

Low dimension ($\dim(\mathbb{B}) \leq 3$)

(i) If $\dim(\mathbb{B}) \leq 2$ and $\mathbb{B}^\tau \neq \emptyset$ then

$$\mathrm{Vec}_{\mathbb{Q}}^{2m}(\mathbb{B}, \tau) \simeq H_{\mathbb{Z}_2}^2(\mathbb{B}|\mathbb{B}^\tau, \mathbb{Z}(1)) .$$

(ii) If $\dim(\mathbb{B}) = 3$ and $\mathbb{B}^\tau \neq \emptyset$ the map

$$\kappa : \mathrm{Vec}_{\mathbb{Q}}^{2m}(\mathbb{B}, \tau) \hookrightarrow H_{\mathbb{Z}_2}^2(\mathbb{B}|\mathbb{B}^\tau, \mathbb{Z}(1))$$

is **only injective** (even though in many situations like $\mathbb{S}^{p,q}$ and $\mathbb{T}^{a,b,c}$ κ turns out to be **bijjective**).

(iii) If $\dim(\mathbb{B}) \leq 3$ and $\mathbb{B}^\tau = \emptyset$ then

$$\mathrm{Vec}_{\mathbb{Q}}^{2m}(\mathbb{B}, \tau) \simeq H_{\mathbb{Z}_2}^2(\mathbb{B}, \mathbb{Z}(1))$$

and

$$\mathrm{Vec}_{\mathbb{Q}}^{2m+1}(\mathbb{B}, \tau) \simeq \begin{cases} H_{\mathbb{Z}_2}^2(\mathbb{B}, \mathbb{Z}(1)) & \text{if } \mathrm{Pic}_{\mathbb{Q}}(\mathbb{B}, \tau) \neq \emptyset \\ \emptyset & \text{if } \mathrm{Pic}_{\mathbb{Q}}(\mathbb{B}, \tau) = \emptyset. \end{cases}$$

Application to involutive spheres $\mathbb{S}^{p,q} := (\mathbb{S}^{p+q-1}, \theta_{p,q})$

$p+q \leq 4$	$q=0$	$q=1$	$q=2$	$q=3$	$q=4$
$\text{Vec}_{\mathbb{Q}}^{2m+1}(\mathbb{S}^{0,q})$	\emptyset	?	?	$2\mathbb{Z} + 1$?
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{S}^{0,q})$	\emptyset	?	?	$2\mathbb{Z}$?
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{S}^{1,q})$	0	0	\mathbb{Z}_2	\mathbb{Z}_2	...
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{S}^{2,q})$	0	?	?	...	
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{S}^{3,q})$	0	?	...		
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{S}^{4,q})$	0	...			

Application to involutive spheres $\mathbb{S}^{p,q} := (\mathbb{S}^{p+q-1}, \theta_{p,q})$

$p+q \leq 4$	$q=0$	$q=1$	$q=2$	$q=3$	$q=4$
$\text{Vec}_{\mathbb{Q}}^{2m+1}(\mathbb{S}^{0,q})$	\emptyset	0	0	$2\mathbb{Z} + 1$	\emptyset
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{S}^{0,q})$	\emptyset	0	0	$2\mathbb{Z}$	0
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{S}^{1,q})$	0	0	\mathbb{Z}_2	\mathbb{Z}_2	...
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{S}^{2,q})$	0	$2\mathbb{Z}$	0	...	
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{S}^{3,q})$	0	0	...		
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{S}^{4,q})$	0	...			

Application to involutive tori $\mathbb{T}^{a,b,c}$ (fixed points)

$$\mathbb{T}^{a,b,c} := \underbrace{\mathbb{S}^{2,0} \times \dots \times \mathbb{S}^{2,0}}_{a\text{-times}} \times \underbrace{\mathbb{S}^{1,1} \times \dots \times \mathbb{S}^{1,1}}_{b\text{-times}} \times \underbrace{\mathbb{S}^{0,2} \times \dots \times \mathbb{S}^{0,2}}_{c\text{-times}}$$

$a + b \leq 3, c = 0$	$a = 0$	$a = 1$	$a = 2$	$a = 3$
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{T}^{a,0,0})$	\emptyset	0	0	0
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{T}^{a,1,0})$	0	$2\mathbb{Z}$?	...
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{T}^{a,2,0})$	\mathbb{Z}_2	?	...	
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{T}^{a,3,0})$	\mathbb{Z}_2^4	...		

Application to involutive tori $\mathbb{T}^{a,b,c}$ (fixed points)

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$a + b \leq 3, c = 0$	$a = 0$	$a = 1$	$a = 2$	$a = 3$
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{T}^{a,0,0})$	\emptyset	0	0	0
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{T}^{a,1,0})$	0	$2\mathbb{Z}$	$(2\mathbb{Z})^2$	\dots
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{T}^{a,2,0})$	\mathbb{Z}_2	$\mathbb{Z}_2 \oplus (2\mathbb{Z})^2$	\dots	
$\text{Vec}_{\mathbb{Q}}^{2m}(\mathbb{T}^{a,3,0})$	\mathbb{Z}_2^4	\dots		

Application to involutive tori $\mathbb{T}^{a,b,c}$ (free involution)

PROPOSITION (D. - Gomi, 2016)

$$\mathbb{T}^{a,b,c} \simeq \mathbb{T}^{a+c-1,b,1} \quad \forall c \geq 2$$

$a + b \leq 2, c = 1$	$a = 0$	$a = 1$	$a = 2$
$\text{Vec}_{\mathbb{Q}}^m(\mathbb{T}^{a,0,1})$	0	?	?
$\text{Vec}_{\mathbb{Q}}^m(\mathbb{T}^{a,1,1})$?	?	...
$\text{Vec}_{\mathbb{Q}}^m(\mathbb{T}^{a,2,1})$?	...	

For all $m \in \mathbb{N}$ **odd** or **even**.

Application to involutive tori $\mathbb{T}^{a,b,c}$ (free involution)

PROPOSITION (D. - Gomi, 2016)

$$\mathbb{T}^{a,b,c} \simeq \mathbb{T}^{a+c-1,b,1} \quad \forall c \geq 2$$

$a + b \leq 2, c = 1$	$a = 0$	$a = 1$	$a = 2$
$\text{Vec}_{\mathbb{Q}}^m(\mathbb{T}^{a,0,1})$	0	\mathbb{Z}_2	\mathbb{Z}_2^2
$\text{Vec}_{\mathbb{Q}}^m(\mathbb{T}^{a,1,1})$	$2\mathbb{Z}$	$\mathbb{Z}_2 \oplus (2\mathbb{Z})^2$...
$\text{Vec}_{\mathbb{Q}}^m(\mathbb{T}^{a,2,1})$	$(2\mathbb{Z})^2$...	

For all $m \in \mathbb{N}$ **odd** or **even**.

Thank you for your attention